

Classical Field Theory NS-364B - Final Exam

The exam will take place on Tue Jul 2 at 15-18 hours in BBG 165/169.

The exam is closed books. 90 points in total; counts as 45% of the total grade. Good luck!

July 1, 2013

1 Theoretical questions (22 points: questions 1-2: 5 pts; 3-4: 6 pts)

1. Sketch the light cone (on which a massless particle can move) and use it to explain the principle of causality. *Hint: it suffices to draw and explain the allowed trajectories. Emphasize what are the properties that trajectories of physical particles must have.*
2. Lorentz symmetry is a non-compact Lie group. Define a Lie group. What is the non-compactness referring to?
3. Electromagnetism possesses a gauge symmetry. Explain how an electromagnetic vector field A^μ transform under the gauge symmetry. Use this gauge transformation to construct the gauge invariant components of A^μ . Explain also which of these components are dynamical and which are non-dynamical. What is the physical meaning of the non-dynamical component(s)?
4. The Liénard-Wiechert potentials are of the form,

$$\varphi(t, \vec{x}) = \frac{e}{R}, \quad \vec{A}(t, \vec{x}) = \frac{\vec{v}}{c} \varphi. \quad (1.1)$$

Explain what physical problem the Liénard-Wiechert potentials solve. In which gauge? Explain what is the meaning of $R = \|\vec{R}\|$ and \vec{v} in (1.1). Explain in detail how R depends on t and \vec{x} . In what sense are these potentials causal?

2 Is this a classical field? (24 points; each question 6 pts; bonus 5 pts)

Consider the following Lagrangian density for two real scalar fields:

$$\mathcal{L} = ic_1 \left[\phi_1 \dot{\phi}_1 + \phi_2 \dot{\phi}_2 - i\phi_2 \dot{\phi}_1 + i\phi_1 \dot{\phi}_2 \right] + \frac{(c_1)^2}{2c_2} [(\nabla\phi_1) \cdot (\nabla\phi_1) + (\nabla\phi_2) \cdot (\nabla\phi_2)] - V(\phi_1^2 + \phi_2^2), \quad (2.1)$$

where V is some potential and the dot represents a time derivative.

1. Calculate the equations of motion for both real fields ϕ_1 and ϕ_2 .
2. Introduce a complex scalar field $\psi = \phi_1 + i\phi_2$. Combine the above two equations of motion appropriately to obtain an equation of motion for ψ and set $c_1 = \hbar$ and $c_2 = m$. What do you observe? - *Hint: You should obtain the time-dependent Schrödinger equation for a single non-relativistic particle.*
3. Now, construct an appropriate Lagrangian density for the complex field ψ ; if you plug in the *Ansatz* $\psi = \phi_1 + i\phi_2$ into your answer, you should retrieve 2.1). Calculate the corresponding Hamiltonian density.
4. Show that the Lagrangian density for the complex scalar field ψ is invariant under the global $U(1)$ transformation $\psi \rightarrow e^{i\alpha}\psi$. Calculate the associated conserved Noether current.
5. *Bonus question:* Comment on why a classical complex scalar field obeys Schrödinger's equation. Do you think quantum mechanics is equivalent to a non-relativistic classical field? Is there any phenomenon in quantum mechanics that you can think of, that you do not expect in Classical Field Theory?

3 Energy momentum tensor (24pts; each question 4pts; bonus 5pts)

Consider a collection of charged particles with position $\vec{x}_n(t)$ and charges e_n . The charge density j^0 and the current density \vec{j} are defined as

$$j^0(\vec{x}, t) = \sum_n e_n \delta^3(\vec{x} - \vec{x}_n(t)), \quad \vec{j}(\vec{x}, t) = \sum_n e_n \dot{\vec{x}}_n(t) \delta^3(\vec{x} - \vec{x}_n(t)) \quad (3.1)$$

The *energy-momentum tensor* of this system is given by

$$T^{\mu\nu} = \sum_n p_n^\mu(t) \dot{x}_n^\nu(t) \delta^3(\vec{x} - \vec{x}_n(t)). \quad (3.2)$$

1. Show that the energy-momentum tensor is only conserved up to a force density G^μ ,

$$\partial_\nu T^{\mu\nu} = G^\mu \quad (3.3)$$

and that G^μ vanishes for free particles (there are no external forces).

2. The electromagnetic force is given by

$$f^\mu \equiv \frac{dp^\mu}{d\tau} = e F^\mu{}_\nu \frac{dx^\nu}{d\tau}, \quad (3.4)$$

where for simplicity we took $e_n = e$. Show that the force density G^μ is given by

$$G^\mu = F^\mu{}_\nu j^\nu, \quad (3.5)$$

where $j^\mu = (j^0, \vec{j})$.

3. To obtain a conserved energy-momentum tensor, we have to include the contribution of the electromagnetic field itself,

$$T_{em}^{\mu\nu} = F^{\mu\rho} F^{\nu\sigma} \eta_{\rho\sigma} - \frac{1}{4} \eta^{\mu\nu} F_{\rho\sigma} F^{\rho\sigma}. \quad (3.6)$$

Show that $T_{tot}^{\mu\nu} = T^{\mu\nu} + T_{em}^{\mu\nu}$ is conserved. *Hint: Use Maxwell's equations.*

4. Show that the total momentum

$$P^\mu = \int d^3x T_{tot}^{\mu 0}(\vec{x}, t) \quad (3.7)$$

is conserved.

Next, we consider the energy momentum tensor of a *perfect fluid*. A comoving (moving along with the fluid) observer will see his/her surroundings as isotropic (this is true only for special configurations of electromagnetic fields, which is what we assume here). In this frame, the particles' energy-momentum tensor (3.2) can be recast into the form

$$T^{\mu\nu} = \begin{pmatrix} \rho & 0 & 0 & 0 \\ 0 & p & 0 & 0 \\ 0 & 0 & p & 0 \\ 0 & 0 & 0 & p \end{pmatrix}, \quad (3.8)$$

where ρ and p are the density and pressure of the fluid.

5. Calculate the components ρ and p of the energy-momentum tensor $T^{\mu\nu}$ for an observer at rest. Assume that the comoving observer has a velocity \vec{v} with respect to the observer at rest.
6. Show that the energy-momentum tensor can be written as

$$T^{\mu\nu} = (p + \rho) \frac{U^\mu U^\nu}{c^2} - p \eta^{\mu\nu}, \quad (3.9)$$

where U^μ is the four-velocity of the fluid, which in the fluid rest frame has the form $U_{rest}^\mu = (c, 0, 0, 0)$.

7. *Bonus question:* By taking the non-relativistic limit of the conservation law $\partial_\mu T^{\mu\nu} = 0$, deduce the following (*Euler's*) *fluid equations*:

$$\partial_t \rho + \nabla \cdot (\rho \vec{v}) = 0, \quad (3.10)$$

$$\partial_t (\rho \vec{v}) + (\vec{v} \cdot \nabla) (\rho \vec{v}) = -\nabla p. \quad (3.11)$$

Of course, these equations hold only when electromagnetic fields are neglected.

Use that, in the non-relativistic limit, the four-velocity is given by $U^\mu = (c, \vec{v})$, where $\|\vec{v}\| \ll c$ and $p \ll \rho$.

4 Advanced Green function (20 points; each question 10 points)

1. The positive and negative frequency Wightman functions for a photon vector field (in Lorentz gauge) or a massless scalar field are defined as homogeneous solutions of the wave equation,

$$-\partial_x^2 G^\pm(x; x') = 0 = -\partial_{x'}^2 G^\pm(x; x'), \quad (4.1)$$

whereby (in the vacuum) $G^+(x; x')$ is determined by the contribution obtained by integrating counterclockwise around the positive frequency pole $k^0 = \omega/c = \|\vec{k}\|$, and $G^-(x; x')$ picks up the contribution by integrating clockwise around the negative frequency pole $k^0 = -\omega/c = -\|\vec{k}\|$. Calculate G^\pm in position space by performing the suitable 4-momentum integrations and show that

$$G^+(x; x') = \frac{-i}{4\pi^2} \frac{1}{\Delta x_+^2}, \quad G^-(x; x') = \frac{-i}{4\pi^2} \frac{1}{\Delta x_-^2}, \quad (4.2)$$

where $i^2 = -1$ and

$$\Delta x_+^2 = (ct - ct' - i\epsilon)^2 - \|\vec{x} - \vec{x}'\|^2, \quad \Delta x_-^2 = (ct - ct' + i\epsilon)^2 - \|\vec{x} - \vec{x}'\|^2. \quad (4.3)$$

Explain the origin of the infinitesimal parameter $\epsilon > 0$ in Eq. (4.3).

2. The Pauli-Jordan, or spectral, two point function can be defined as

$$G_{PJ} = G^- - G^+. \quad (4.4)$$

Show that

$$G_{PJ} = -\frac{\text{sign}(t-t')}{2\pi} \delta(\Delta x^2 - \epsilon^2), \quad (4.5)$$

where Δx^2 is defined as $\Delta x^2 = c^2(t-t')^2 - \|\vec{x} - \vec{x}'\|^2$, $\text{sign}(t-t') = \Theta(t-t') - \Theta(t'-t)$ and $\Theta(t-t')$ denotes the Heaviside function. Calculate the advanced Green function in position space,

$$G_a(x; x') = \Theta(t' - t) G_{PJ}(x; x'). \quad (4.6)$$

Explain the causality structure of the advanced Green function. Explain also what kind of problems G_a can be used to solve.

Hint: Make use of the Plemelj-Sokhotski theorem,

$$\frac{1}{x \mp i\epsilon} = \mathcal{P} \frac{1}{x} \pm i\pi \delta(x),$$

where \mathcal{P} denotes a principal value (when integrating) and $\epsilon > 0$ is an infinitesimal parameter.